

GRAVITY EOSC. 450

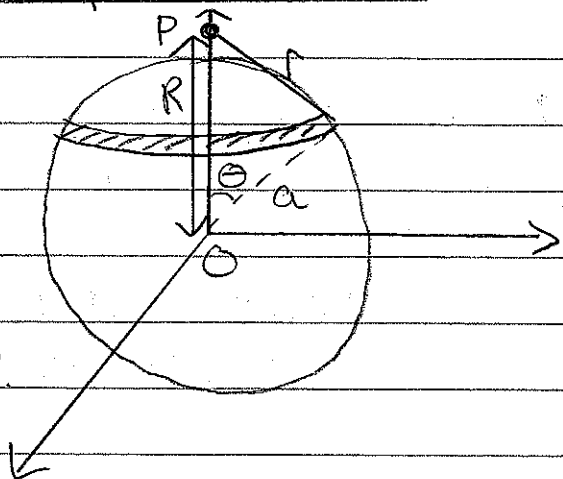
FORWARD MODELING: Analytical Examples

See also Blakely pages 49-59

Example 1: Gravitational Potential + Gravity Field due to a Spherical Shell

Blakely  
p. 49-51

- 1) Shell, radius =  $a$
- 2) Pt. P outside spherical shell
- 3) Origin @ center of shell.
- 4) For convenience, orient coord system to go thru' P.



- 1] Procedure :
- a) calculate pot<sup>l</sup>  $\Phi$  from  $\Phi = G \int \frac{\rho(r)}{r} dv$
  - b) calc.  $\vec{g}$  from  $\vec{g} = -\vec{\nabla} \Phi$

2] Assume mass concentrated on spherical shell so

$$\int \rho dv \rightarrow \int \sigma ds$$

$\rightarrow$  surface mass density  
 = mass per unit area.

$$\text{so } \Phi(P) = G \int_s \frac{\sigma(s) ds}{r}$$

$\rightarrow ds = a^2 \sin\theta d\theta d\phi$

$$\Phi(P) = G \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \frac{\sigma a^2 \sin\theta d\theta d\phi}{r}$$

$$\Phi(P) = G\sigma^2 \int_0^{2\pi} \int_0^{\pi} \frac{\sin\theta}{r} d\theta d\phi \quad (2)$$

$$= G\sigma^2 2\pi \int_0^{\pi} \frac{\sin\theta}{r} dr$$

Issue:  $r$  is a function of  $\theta$  in our picture.

3 Use Cosine Law

$$r^2 = R^2 + a^2 - 2aR \cos\theta$$

$$\Rightarrow 2r dr = -2aR (-\sin\theta) d\theta$$

$$\Rightarrow \frac{\sin\theta}{r} d\theta = \frac{dr}{aR}$$

to allow

→ change of variables:

$$\theta = 0 \Rightarrow r^2 = R^2 + a^2 - 2aR = (R-a)^2$$

$$r = \pm (R-a)$$

We'll use  $r = R-a$  [WHY] ??

$$\theta = \pi \Rightarrow r^2 = R^2 + a^2 + 2aR = (R+a)^2$$

$$r = \pm (R+a)$$

We'll use  $r = R+a$  [WHY] ??

$$4 \text{ So } \Phi(P) = 2\pi a^2 G \sigma \int_{r_1}^{r_2} \frac{dr}{aR} = 2\pi a^2 G \sigma \int_{R-a}^{R+a} \frac{dr}{aR}$$

$$= 2\pi a^2 G \sigma \left[ \frac{(R+a) - (R-a)}{aR} \right]$$

$$= \frac{4\pi a^2 G \sigma}{R} = \frac{GM}{R}$$

$$\Rightarrow \Phi(P) = \frac{GM}{r}$$

## Field Modeling (cont) (Spherical shell)

(3)

- As there is rotational symmetry, we could have chosen  $P$  in any dir<sup>n</sup> & arrived @ same answer.
- Note: Grav<sup>l</sup> pot<sup>l</sup> @ point,  $P$ , outside a spherical shell of surface mass density  $\sigma$  is the same as if the mass of the shell were all concentrated @ the center.

### 5 More generally (outside shell)

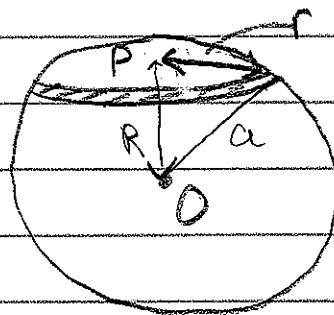
$$\Phi(\vec{r}) = \frac{GM}{r} \quad \vec{g}(\vec{r}) = -\frac{GM}{r^2} \hat{r}$$

since  $\vec{g} = -\nabla\Phi$

Remark: can obtain same result using Gauss's thm.

### 6 Inside the shell

- Point  $P$  inside shell  
dist.  $R$  from origin
- shell radius  $a$ .



$$\Phi(P) = \frac{GM}{a}$$

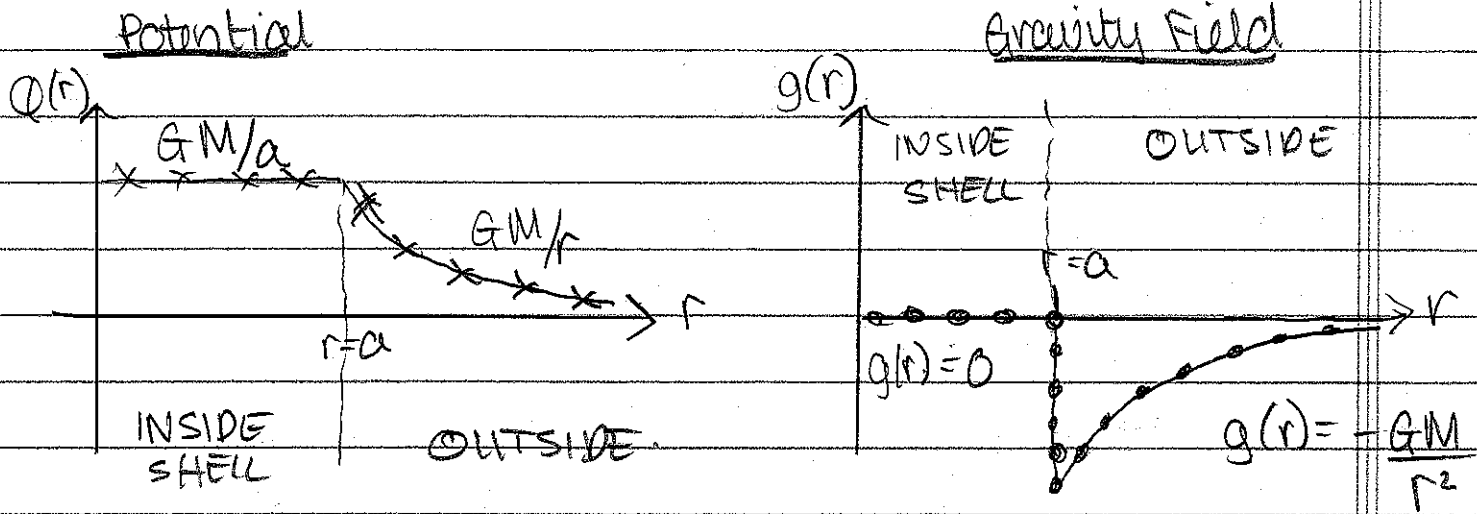
STUD THIS  $\uparrow$

$$\vec{g} = \nabla\Phi = 0$$

and so  $\nabla^2\Phi = 0$

# Spherical shell

(4)



Potential due to solid sphere

- see next 2 pages ; also Blehely p. 51-54

# Fluid Modeling (ctd)

(5)

## Example: Solid Sphere

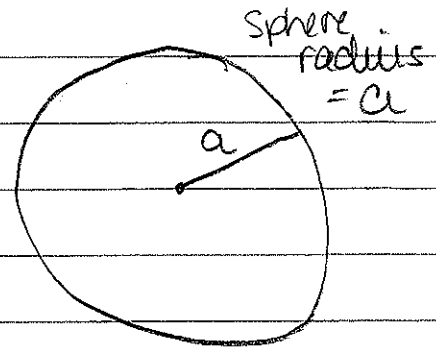
- follows from results for shell + superposition

### A) Outside sphere

$$\Phi(r) = \frac{GM}{r}, \quad r > a$$

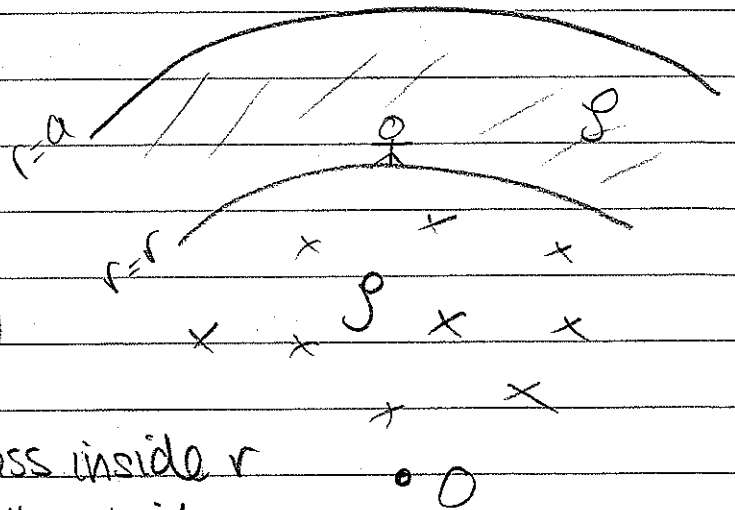
and  $\nabla^2 \Phi = 0$

$$M = \sum_{i=1}^N m_i^{\text{shell}} = \frac{4\pi}{3} \bar{\rho} a^3$$



### B) Inside sphere

Imagine an observer @ radius 'r' inside the planet



The pot<sup>l</sup> @ r can be written

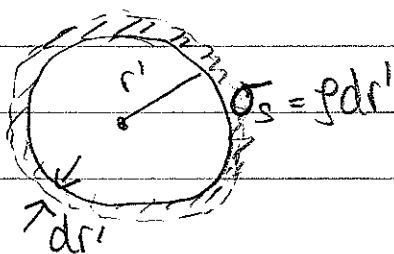
$$\Phi(r) = \Phi_o(r) + \Phi_I(r)$$

$\Phi_I(r)$  pot<sup>l</sup> @ r due to mass inside r

$\Phi_o(r)$  " " " " " " outside r

$$\Phi_I(r) = \frac{GM}{r} = \frac{G}{r} \frac{4\pi r^3 \rho}{3} = \frac{4}{3} \pi r^2 G \rho$$

$$\Phi_o(r) = \int_r^a \left( \text{pot}^l \text{ due to a spherical shell @ } r' \right) dr' \quad r \leq r' \leq a$$



$$\begin{aligned} \Phi(r < r') &= \frac{GM}{r'} = \frac{G(4\pi \sigma_s r'^2)}{r'} \\ &= 4\pi G \sigma_s r' = 4\pi G \rho r' dr' \end{aligned}$$

$$\text{so } \phi_0(r) = \int_r^a 4\pi G \rho r' dr' = 2\pi G \rho (a^2 - r^2) \quad (6)$$

$$\begin{aligned} \text{so } \phi(r) &= \phi_i(r) + \phi_0(r) \\ &= \frac{4}{3}\pi r^2 G \rho + 2\pi G \rho (a^2 - r^2) \end{aligned}$$

$$\Rightarrow \boxed{\phi(r) = \frac{2\pi G \rho}{3} (3a^2 - r^2)}$$

Potential inside solid sphere of uniform density

$$\vec{g} = -\nabla \phi(\vec{r}) = \hat{r} \frac{\partial}{\partial r} \phi(r) = -\frac{4}{3}\pi G \rho r \hat{r}$$

$$\Rightarrow \boxed{\vec{g} = -\frac{4\pi G \rho r}{3} \hat{r}}$$

i.e. gravity varies linearly w/  $r$  inside the sphere.

SHOW THAT POISSON'S EQ<sup>n</sup> HOLDS INSIDE THE SPHERE

$$\text{i.e. } \boxed{\nabla^2 \phi = -4\pi G \rho}$$