

# Brittle strain regime transition in the Afar depression: Implications for fault growth and seafloor spreading

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## ABSTRACT

**Numerical and analogue models of fault-population evolution suggest that we should see a frequency-size transition from power law to exponential as brittle strain within a region increases. We now observe this transition in continental rift faults in the Afar region, Ethiopia and Djibouti. Furthermore, we observe that this strain transition is accompanied by a plateau in fault density and an increase in the displacement:length ratio of faults with increasing brittle strain. Once faults reach a critical density, stress shadows of nearby faults inhibit fault nucleation and restrict tip propagation. However, pinned faults continue to accommodate extension by accumulating more displacement, thus leading to higher displacement:length ratios in the higher strain regime. This strain regime transition is important to understanding the physics of fault growth and may also be important to rift localization and the development of plate boundaries.**

**Keywords:** Afar, fault scaling, brittle strain, fault interaction.

## INTRODUCTION

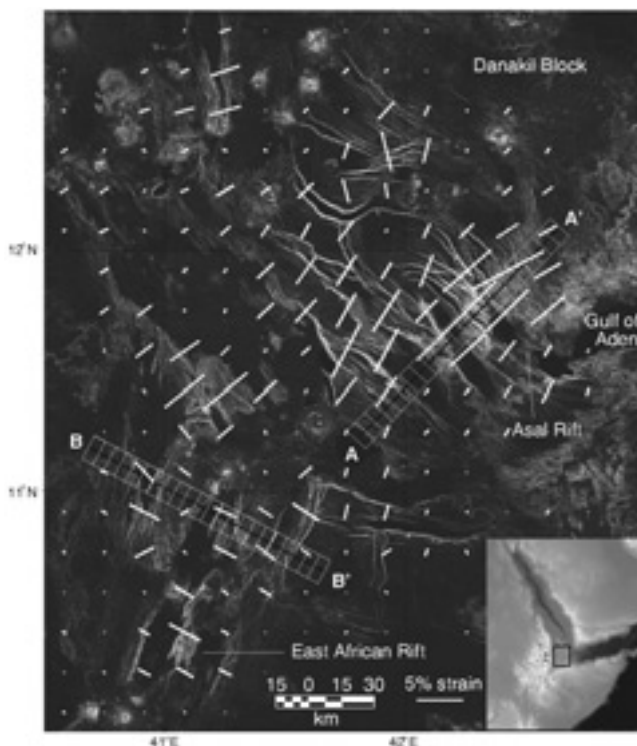
Observations of fault size, specifically displacement and length distributions, provide important constraints on physical models of fault growth as well as tests of model predictions. Numerical and clay models of fault-population evolution show that the size distribution of cracks has a transition from power law to exponential with increasing strain (Spyropoulos, 1999; Spyropoulos et al., 1999; Ackermann et al., 1997), and that crack density has a maximum value (Spyropoulos et al., 1999). Most observations of fault populations growing in low-strain regions show power law frequency-length scaling (Scholz and Cowie, 1990; Scholz et al., 1993; Scholz, 1997; Bohnenstiehl and Kleinrock, 1999). In higher strain settings, such as along the flanks of some mid-ocean ridges, faults have been observed to follow exponential frequency-size distributions (Cowie et al., 1993; Carbotte and Macdonald, 1994). Although Spyropoulos et al. (1999) suggested that these observations represent the end-member cases of their model observations, nowhere has this transition been observed in a single tectonic setting.

Another important scaling characteristic of fault populations is the linear relationship between fault displacement and length (Cowie and Scholz, 1992a; Dawers et al., 1993; Schlische et al., 1996). However, large variations in displacement:length ratios have been observed (Walsh and Watterson, 1988; Marrett and Allmendinger, 1991; Cowie and Scholz, 1992b; Gillespie et al., 1992). Much of the variation within similar tectonic settings has been attributed to fault interaction (Dawers and Anders, 1995; Cartwright et al., 1995; Wojtal, 1996; Willemse et al., 1996). Interaction can account for at least half the scatter in displacement:length

ratios (Cartwright et al., 1995; Willemse et al., 1996; Gupta and Scholz, 2000). Poulimenos (2000) observed that higher displacement:length ratios correlate with high brittle strain regions in the Western Corinth graben, Greece. He is the first to link increasing strain to changes in displacement:length ratios. We see that for faults from the Afar triangle, not only are displacement:length ratios related to strain, but changes in fault size distribution, displacement:length ratios, fault density, and brittle strain are all intimately related to one another.

## GEOLOGIC SETTING

The triple junction between the Arabian, Nubian, and Somalian plates is shown in Figure 1 (inset). Motion is accommodated by the East African Rift, the Sheba Ridge, the Red Sea Ridge, and faults in the Afar triangle. The three major plate boundaries are in different stages of evolution. Around the Sheba Ridge in the Gulf of Aden clear magnetic anomalies are identified to 10 Ma, although spreading may have started earlier (Cochran, 1981). The Red Sea began spreading 4–5 Ma (Roeser, 1975; Searle and Ross,



**Figure 1. Location map and extension vectors. Inset shows extensional triple junction. Arabian plate is to north, Nubia is to west, and Somalia is to east. Area of larger image is outlined in inset. Larger image shows slopes calculated from digital elevation model and average extension vectors (thick white lines) for  $18.5 \times 18.5$  km<sup>2</sup> regions. Maximum average extension for 342 km<sup>2</sup> regions is 9.0%. Steep slopes appear brighter. Linear features are faults. Circular features, such as those in northwest, are volcanoes. Fluvial channels are apparent on eastern edge of image. Transects mark locations of detailed fault-population analyses.**

1975), while the East African Rift remains a continental rift.

The extended crust in the Afar triangle is influenced by all three plate boundaries as well as the Danakil block. The Danakil block is thought to have started as a piece of the Nubian plate and rotated counterclockwise  $23^\circ$  to its current position (Sichler, 1980; Collet et al., 2000). The relatively undeformed Danakil block forms the northern boundary of the Afar triangle. To the southwest and southeast, the Ethiopian Plateau looms 2 km above the floor of the Afar triangle. The plateau consists of Tertiary basalt formed through the influence of the Afar mantle plume (e.g., Mohr, 1983; Hofmann et al., 1997).

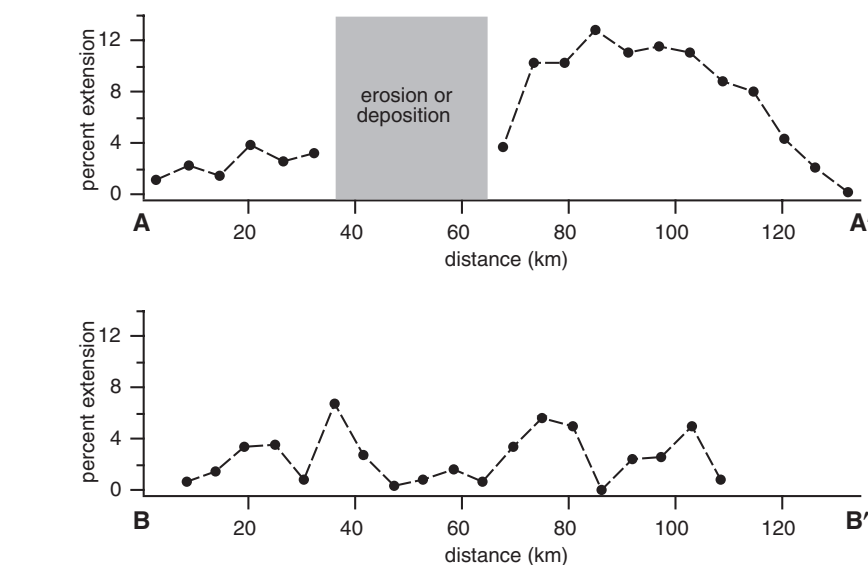
The nature of the crust in the Afar triangle is controversial (Makris and Ginsburg, 1987; Mohr, 1989). The crust is of intermediate thickness (15–25 km), and shows a velocity structure transitional between continental and oceanic (Berckhemer et al., 1975). The strength of the crust, measured as effective elastic thickness, decreases from 17 to 5 km in highly extended portions of Afar (Ebinger and Hayward, 1996). Much of the Afar triangle is covered in Pliocene basalt of the Stratoid series (Barberi and Varet, 1977). The northwest-southeast-trending Asal rift (Fig. 1) on the eastern edge of the Afar triangle is widely cited as a location where the beginnings of seafloor spreading can be observed on land (e.g., Stein et al., 1991).

We studied in detail the extension and fault-population characteristics across two transects. Faults that cross transect A–A' are thought to be due to extension along the Arabia-Somalia plate boundary (Manighetti et al., 1998). However, the region just north of A' belongs to the Danakil block, not Arabia (e.g., Collet et al., 2000). In any case, these faults are approximately perpendicular to the Arabia-Somalia motion vector (e.g., Jestin et al., 1994). The Asal rift (Fig. 1) is thought to accommodate most of the predicted present-day motion along this plate boundary (De Chabaliere and Avouac, 1994). Faults that cross transect B–B' clearly belong to the northernmost portion of the East African Rift, the boundary between the Nubian plate to the west and the Somalian plate to the east.

## STRAIN REGIME TRANSITION

### Regional Extension

We measured extension for the area shown in Figure 1. Surface extension due to normal faults was measured using a 3 arcsecond digital elevation model of the region that provides 92 m horizontal and 10 m vertical resolution. Because the region has been recently resurfaced by volcanics, most of the high-frequency variability in the topography can be attributed to normal faulting. We calculated the vertical elevation difference between adjacent pixels. When this difference was not due to an obvious volcanic, fluvial, or erosional construct, we assumed that the eleva-



**Figure 2.** Average extension for  $6 \times 12 \text{ km}^2$  regions along A–A' and B–B' (indicated by white rectangles in Fig. 1). Maximum average extension is 12.9% for  $72 \text{ km}^2$  subregions along A–A' and B–B'. Assuming 30% increase in fault density with uniform offset of 10 m, estimated error in extension plots within symbol size.

tion difference was created by normal faulting. Several weeks of field reconnaissance in the Asal rift confirmed the predominance of normal faulting. Observations of faults in the Asal rift indicate that they typically dip steeply,  $\sim 70^\circ$ . Because faults may become steeper as they approach the surface, we assumed a crustal average fault dip of  $60^\circ$  for our calculation. Given fault throw and dip, we calculated the resulting heave or extension using simple trigonometry. Thus the extension estimates will vary with the cotangent of the fault dip. By assuming fault dip rather than using slopes calculated from the digital elevation model, we also minimized underestimation of fault dip due to undersampling or slumping. Changes in topography due to minor strike-slip motion, buried or eroded offset, deformation below data resolution, and low-frequency topographic variation are neglected in this analysis. Hence, the extension values presented in this study can be viewed as minimum values. We do not calculate the entire strain tensor.

Extension measured in the Afar triangle is shown in Figure 1. The thick white lines indicate the magnitude and direction of extension averaged in  $18.5 \text{ km}$  by  $18.5 \text{ km}$  regions. The extension averaged in this size region ranged from 0.0% to 9.0%. The primary features are the change in extension direction from northeast across the Asal rift to east-southeast across the East African Rift and the broad regions of strain localization centered on those rifts.

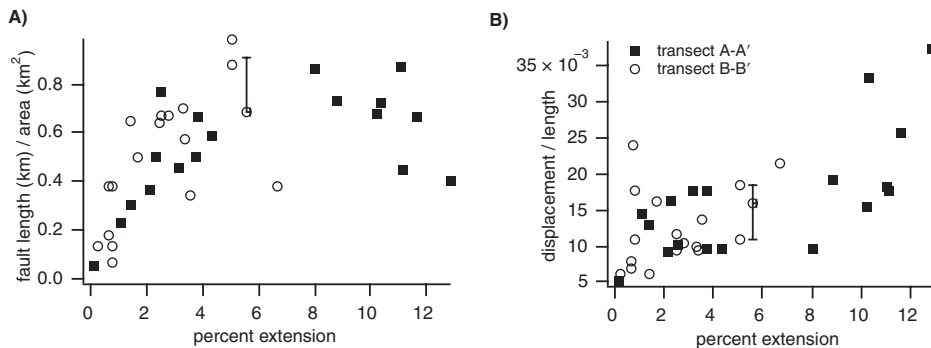
Detailed fault-population analysis was conducted along transects A–A' and B–B'. Extension averaged in smaller subregions along these transects (indicated by white rectangles in Fig. 1) is shown in Figure 2. The zone of extension along both transects is about 100–120 km wide. In tran-

sect A–A' maximum extensional strain within subregions is at about the 85 km point, just to the northwest of the Asal rift. Here it is believed that the beginning stages of slow seafloor spreading can be observed on land (Stein et al., 1991; De Chabaliere and Avouac, 1994; Manighetti et al., 1998). We measure total brittle extension across transect A–A' to be 8.5 km. In contrast, the extension across the East African Rift, transect B–B', is diffuse and there is not one well defined extensional peak. The total extension across this transect, where seafloor spreading is not yet imminent, is 3.4 km.

### Fault Scaling

Displacement length profiles for normal faults that are partially or fully contained within subregions shown in Figure 1 were extracted from the digital elevation model. Transects along the hanging wall were subtracted from profiles along the footwall to obtain displacement and length information. Fault density is calculated as fault length per unit area ( $\text{km}/\text{km}^2$ ) (Fig. 3). We estimate that fault lengths may be underestimated by 0.4–4.0 km, depending on the displacement gradients near tips. Lengths are underestimated because we lose resolution below 10–15 m of vertical offset. The error will be more significant for small faults because the missing length will represent a larger fraction of their total length. We have added a constant value of 2 km to fault lengths in Figure 4 to help account for the missing length (assuming that vertical offset below 12 m is not observed, and an observed displacement gradient of 0.012 near the tips).

We found that the fault density and displacement:length ratios change with the estimated brittle strain. At first the fault density increases ap-



**Figure 3. A:** Fault density (kilometers of fault length per unit area) as function of estimated strain. Example error bar shows effect of 16 km of additional fault length with 10 m uniform offset. Extension estimate does not change significantly. **B:** Average displacement:length ratios for faults that are entirely or partially contained within subregions along transects plotted against percent extension. Example error bar shows uncertainty in displacement:length ratios and percent extension due to digital elevation model resolution. It is more likely that displacement:length ratios are slightly overestimated than underestimated because maximum offsets are accurate to within 10 m, but fault length could be underestimated by 0.4–4 km.

proximately linearly with strain, but at ~6%–8% the density levels off and does not increase further (Fig. 3A). Note that until about 8% extension, the displacement:length ratio remains relatively constant around 0.012 (Fig. 3B). Beyond 8% strain displacement:length ratios rise sharply from a mean value of 0.012 to 0.024. Brittle strain,  $\epsilon$ , is proportional to  $LWD$ , where the width  $W$  is a constant width,  $L$  is fault length, and  $D$  is displacement. In the low-strain regime strain accumulates with increasing fault density (total fault length), whereas in the high-strain regime strain accumulates because of increased displacement on preexisting faults.

We also observe an evolution in the fault-population statistics (Fig. 4). The faults in subregions that accommodate <8% extension (filled circles) exhibit close to a power law frequency length distribution above the resolution cutoff at 4–5 km (Fig. 4A). In contrast, the fault size distribution in subregions that accommodate >8% extension (squares) is closer to an exponential

(Fig. 4B). The maximum displacement distribution (not shown) exhibits the same basic relationships as the fault length distribution. Faults in low-extension subregions show power law scaling and faults in high-extension subregions show exponential scaling.

#### INTERPRETATION

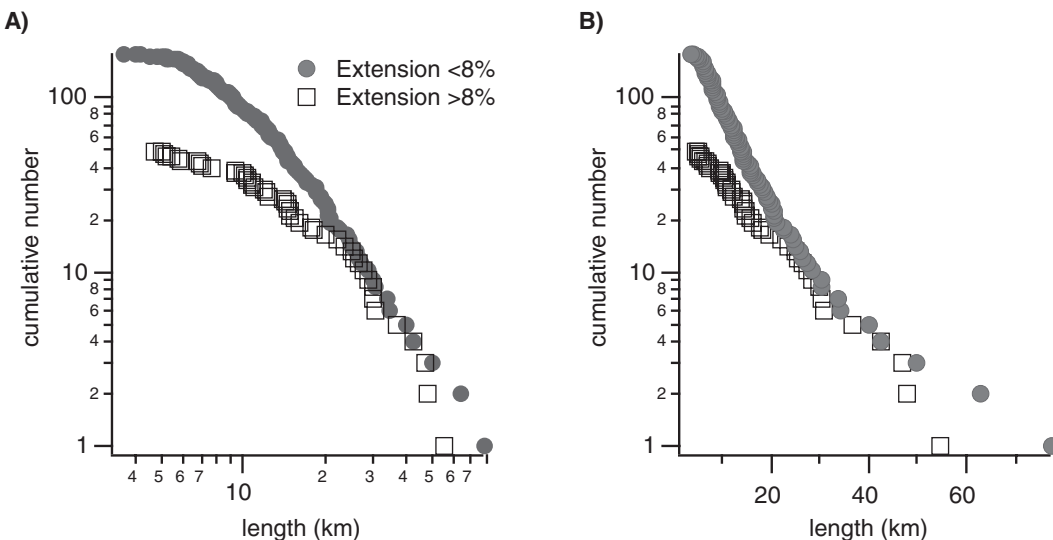
Results show that brittle strain initially accumulates as a result of steady increase in fault density, but that the rate of increase in fault density with strain gradually diminishes until it reaches a steady-state value at 6%–8% strain, beyond which further strain is produced by increasing displacement:length ratios. The fault size distribution is power law in the low-strain regime but becomes exponential at high strain. Higher strain regions have relatively fewer small faults than regions with lower strain. These observations are the same as those observed by Spyropoulos et al. (1999) in a clay model. They interpreted these phenomena as resulting from an increase in fault

interaction, nucleation of new faults being inhibited by stress shielding around preexisting faults and older faults being annihilated by coalescence (see also Ackermann and Schlische, 1997; Cowie, 1998). We have found, in addition, that high strains are accommodated by a progressive increase in the displacement:length ratio, which implies that, statistically, displacement is accommodated on faults that are no longer growing in length. This occurs because of pinning of fault growth by stress interaction with other faults (Gupta and Scholz, 2000; Contreras et al., 2000).

Poulimenos (2000) also observed higher displacement:length ratios in normal faults that occurred in higher strain settings in Greece. However, both high- and low-strain populations from the western Corinth graben showed power law frequency-length scaling. We favor Poulimenos' second interpretation for these observations. The faults in the higher strain region trend west-northwest and are crosscut by a population of smaller north-northeast-trending faults. These smaller faults serve to pin the west-northwest-trending fault population as stress fields do in Afar, preventing tip growth. The frequency-length distributions suggest that the west-northwest-trending faults from Greece have not entered a stage of growth dominated by coalescence.

#### CONCLUSIONS

We have observed fault populations in a single rock type and tectonic setting in different stages of evolution in the Afar triangle, a scenario previously only modeled. A strain regime transition from power law to exponential frequency-size scaling occurs when faults reach a certain density (about 0.6 km of fault length per square kilometer for the faults in Afar). Once this critical density is reached, little new fault surface area is created. Instead, faults must lengthen primarily by coalescence rather than by growth or nucleation. Saturation of fault surface area appears to occur as faults become pinned by the stress fields of



**Figure 4. A:** Cumulative number of faults vs. fault length, plotted on log-log axes. Cumulative number is calculated as number of faults with length greater than or equal to particular value. Power law relationships plot as straight line. **B:** As in A, except data are plotted on log-linear axes. Exponential relationships plot as straight line.

nearby faults. Their tips cannot easily propagate into stress shadows around nearby faults; new faults are also inhibited from nucleation by stress shadows. Displacement:length ratios must increase to balance stress shadows from nearby faults against stress concentrations near tips. In Afar, fault interaction plays a pivotal role in the strain transition.

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